

## **ELI-NP at Magurele - “Pulse and Impulse of ELI”**

- 1)"**Polaritonic pulse** and coherent X- and gamma rays from Compton (Thomson) backscattering" (MApostol&MGanciu), J. Appl. Phys. **109** 013307 (2011) (1-6)
- 2)"Dynamics of **electron–positron pairs** in a vacuum polarized by an external radiation field" (MA), Journal of Modern Optics, **58** 611 (2011)
- 3)"**Classical interaction** of the electromagnetic radiation with two-level polarizable matter" (MA), Optik **123** 193 (2012)
- 4)"**Coherent polarization** driven by external electromagnetic fields" (MA&MG), Physics Letters **A374** 4848 (2010)

- 5)"Coupling of **(ultra-) relativistic atomic nuclei** with photons" (MA&MG), AIP Advances **3** 112133 (2013)
- 6)"Propagation of **electromagnetic pulses** through the surface of dispersive bodies" (MA), Roum J. Phys. **58** 1298 (2013)
- 7)"**Giant dipole oscillations** and ionization of heavy atoms by intense electromagnetic pulses" (MA)...

## **What we should know firstly about lasers:**

**Intensity**  $I \simeq \mathcal{E}/d^2\tau$ ;

- 1) Threshold (1990) :  $I = 10^{14} - 10^{15} w/cm^2$  (CPA);  $E \simeq \sqrt{I/c} \simeq 10^6 (esu)(3 \times 10^4 V/m)$  - atomic fields - ionization, HO harmonics
- 2) Current:  $I = 10^{18} - 10^{19} w/cm^2$ ,  $E \simeq 10^8$  - relativistic, non-linear
- 3) Few claims:  $10^{21} - 10^{22} w/cm^2$  (ultrahigh),  $E \simeq 10^9$  - multi-photon, radiation reaction, q vacuum, particle creation

Schwinger limit  $E \simeq 10^{13}$

**Power**  $P \simeq \mathcal{E}/\tau$  ( $Id^2$ )

Current:  $\tau \simeq x \cdot 10^{-15}s$ ,  $d \simeq 3x \cdot 10^{-5}cm$

**The petawatt:**  $10^{21} \times (9x^2)10^{-10} = 10^{15}w$  ( $x = 50$ )!

**Wavelength**  $\varepsilon = 1eV$ ,  $\omega = 10^{15}s^{-1}$ ,  $\lambda = 10^{-5}cm$  (tens of  $\lambda$  in a pulse)

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**Giant dipole oscillations and ionization of heavy atoms by  
intense electromagnetic fields**

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## Thomas-Fermi model (1927)

$$\hbar^2 k_F^2 / 2m - e\varphi = 0 , \quad k_F = (2me/\hbar^2)^{1/2} \varphi^{1/2}$$

Poisson's equation

$$\Delta\varphi = -4\pi Ze\delta(\mathbf{r}) + 4\pi ne , \quad n = k_F^3 / 3\pi^2 = (1/3\pi^2)(2me/\hbar^2)^{3/2} \varphi^{3/2}$$

$$\Delta\varphi = -4\pi Ze\delta(\mathbf{r}) + \frac{4e}{3\pi}(2me/\hbar^2)^{3/2} \varphi^{3/2}$$

**Solution:** Numerical

Binding energy  $E \simeq -20.8Z^{7/3}$ eV, exact for  $Z \rightarrow \infty$  (Lieb&Simon, 1977)

Empirical:  $E \simeq -16Z^{7/3}$ eV ( $Z > 20$ )

Corrections: (asymptotic) series in  $Z^{-1/3}$ ; (convergence?)

Another drawback: no-binding theorem (Teller, 1962).

## Linearized version (with L C Cune; binding)

$k_F^2 \rightarrow 2\bar{k}_F k_F$ ,  $k_F^3 \rightarrow 3\bar{k}_F^2 k_F$ , etc (quasi-classical approximation)

Alternatively,

$$\frac{\hbar^2}{m} k_F \delta k_F - e \delta \varphi = 0 , \quad k_F = (me/\hbar^2 \bar{k}_F) \varphi$$

$$n = (me\bar{k}_F/\pi^2 \hbar^2) \varphi = (q^2/4\pi e) \varphi$$

Thomas-Fermi screening wavevector  $q$

$$q^2 = 4me^2 \bar{k}_F / \pi \hbar^2 = \frac{4\bar{k}_F}{\pi a_H} , \quad a_H = \frac{\hbar^2}{me^2} = 0.53 \text{Å}$$

$$\Delta \varphi = -4\pi Z e \delta(\mathbf{r}) + q^2 \varphi , \quad \varphi = \frac{Ze}{r} e^{-qr}$$

## Energy

$$E_{kin} = V \hbar^2 k_F^5 / 10\pi^2 m \rightarrow (\hbar^2 \bar{k}_F^4 / 2\pi^2 m) \int d\mathbf{r} \cdot \mathbf{k}_F = \frac{\pi e a_H^3}{128} q^6 \int d\mathbf{r} \cdot \boldsymbol{\varphi}$$

$$E_{kin} = \frac{\pi^2 a_H^3}{32} Z e^2 q^4$$

$$\begin{aligned} E_{pot} &= \int d\mathbf{r} (\rho_e \varphi - \frac{1}{2} \rho_e \varphi_e) = \frac{1}{2} \int d\mathbf{r} (\rho_e \varphi + \rho_e \varphi_c) = \\ &= -\frac{e}{2} \int d\mathbf{r} n(\varphi + \varphi_c) = -\frac{q^2}{8\pi} \int d\mathbf{r} (\varphi^2 + \varphi \varphi_c) \\ \rho_e &= -en, \quad \varphi_e = \varphi - \varphi_c, \quad \varphi_c = Ze/r \end{aligned}$$

$$E_{pot} = -\frac{3}{4} Z^2 e^2 q$$

## Total energy

$$E = E_{kin} + E_{pot} = \frac{\pi^2 a_H^3}{32} Z e^2 q^4 - \frac{3}{4} Z^2 e^2 q$$

Minimum value

$$E = -\frac{9 \cdot 6^{1/3}}{16\pi^{2/3}} Z^{7/3} \frac{e^2}{a_H} = -0.42 Z^{7/3} \frac{e^2}{a_H} = -11.4 Z^{7/3} eV$$

for

$$q = (6/\pi^2)^{1/3} \frac{Z^{1/3}}{a_H} = 0.85 \frac{Z^{1/3}}{a_H}$$

(atomic unit for energy  $e^2/a_H \simeq 27.2 \text{eV}$ )

## **Electron distribution:**

$r^2 n$ , maximum for  $R \sim 1/q \sim Z^{-1/3} a_H$ , the "radius" of the electronic charge

Quasi-classical description:  $a_H/Z \ll R \sim a_H/Z^{1/3} \ll a_H$

$(a_H = \hbar^2/m e^2 = 0.53\text{\AA})$

## Quantum correction

$$-\frac{e}{v} \int_v d\mathbf{r} \cdot \varphi = -3Ze^2 q \cdot \frac{1}{x^3} (1 - e^{-x} - xe^{-x}) , \quad v = 4\pi R^3/3 , \quad x = qR$$

Multiply by  $\int_v d\mathbf{r} \cdot n$

$$\Delta E = \frac{16}{3} E \cdot \frac{1}{x^3} (1 - e^{-x} - xe^{-x})^2$$

Minimize with respect to  $R$  (maximum 0.073 for  $x \simeq 0.75$ ),

$$\Delta E = 0.39E = -4.44Z^{7/3}\text{eV}$$

Total energy

$$E = -11.4Z^{7/3}\text{eV} - 4.44Z^{7/3}\text{eV} = -15.84Z^{7/3}\text{eV}$$

Compare with empirical  $E \simeq -16Z^{7/3}\text{eV}$

## Giant dipole oscillations

$$Rq = \text{const}; u = -(1/q^2)\delta q$$

$$\delta E = \delta E_{kin} = \frac{27}{4\pi^2} \frac{Z^3 e^2}{a_H^3} u^2$$

frequency  $\omega_0$ ,  $\delta E = \frac{1}{2}M\omega_0^2 u^2$ ,  $M = Zm$

$$\omega_0 = \left( \frac{27}{2\pi^2} \right)^{1/2} \frac{Ze}{\sqrt{ma_H^3}} \simeq 4.5Z \times 10^{16} s^{-1}$$

energy  $\hbar\omega_0 \simeq 28ZeV$ , moderate  $X$ -rays

Wavelength  $\lambda_0 = 2\pi c/\omega_0 \simeq \frac{4.2}{Z} \times 10^{-6} cm$  much larger than the dimension of the atom

**Damping force:**  $F_d = 2Q^2\ddot{v}/3c^3$ ,  $Q = -Ze$ ,  $v = \dot{u}$

$$\ddot{v} = \omega_0^2 v, F_d = \frac{2Z^2 e^2}{3c^3} \omega_0^2 v = M\gamma \dot{u}, M = Zm$$

$$\gamma = \frac{2Ze^2}{3mc^3} \omega_0^2$$

Note:  $\gamma \ll \omega_0$ , since

$$\frac{2}{3}Z \frac{e^2}{mc^2} \ll \frac{c}{\omega_0}, \quad \frac{4\pi}{3} Z r_0 \ll \lambda_0$$

Class em radius  $e^2/mc^2 = r_0 \simeq 2.8 \times 10^{-13} \text{ cm}$

Quality ratio (natural breadth of the spectroscopic line)

$$\frac{\gamma}{\omega_0} = \frac{4\pi Z r_0}{3\lambda_0} \simeq 2.8 Z^2 \times 10^{-7}$$

## Equation of motion

$$M\ddot{u} + M\omega_0^2 u + M\gamma\dot{u} = QE \cos \omega t, \quad m\ddot{u} + m\omega_0^2 u + m\gamma\dot{u} = -eE \cos \omega t$$

Solution

$$u = a \cos \omega t + b \sin \omega t$$

$$a = \frac{eE}{m} \frac{\omega^2 - \omega_0^2}{(\omega^2 - \omega_0^2)^2 + \omega^2\gamma^2}, \quad b = -\frac{eE}{m} \frac{\omega\gamma}{(\omega^2 - \omega_0^2)^2 + \omega^2\gamma^2}$$

Energy conservation, power loss

$$P = m\gamma\dot{u}^2 = \frac{e^2 E^2}{2m} \frac{\gamma\omega^2}{(\omega^2 - \omega_0^2)^2 + \omega^2\gamma^2}$$

## Resonance

$$P_{res} = \frac{e^2 E^2}{2m\gamma}$$

Moderate fields  $E = 1/300 esu$  ( $100 V/m$ ):

$$P_{res} = \frac{1}{Z^3} \times 10^{-7} erg/s$$

Transition rate

$$R = P_{res}/\hbar\omega_0 \simeq \frac{2}{Z^4} \times 10^3 s^{-1}$$

(Radiation)

## Partial ionization

Linear Th-F model:

$$\omega'_0 = (\delta Z/Z)\omega_0$$

$$\gamma' = (2e^2\delta Z/3mc^3)\omega_0'^2$$

(UV)

## Anharmonicities and ionization

Amplitude (at resonance) (compare with  $a_H = 0.53\text{\AA}$ )

$$|b_0| = \frac{eE}{m\omega_0\gamma} = \frac{8}{Z^4} \times 10^{-10} E \text{ cm}$$

Small, harmonic approximation; higher-order corrections

Potential energy ( $q = 0.85Z^{1/3}/a_H$ )

$$U = \frac{1}{2}M\omega_0^2 u^2 \left( 1 - \frac{8}{3}qu + \frac{31}{6}q^2 u^2 + \dots \right)$$

Combined frequencies, harmonics, freq shifts

Successive approximations, SCHOA

**Ionization**  $((q + \delta q)(R + u) = 1 \ (\delta q = -q^2 u / (1 + qu))$

$$\Rightarrow U_\infty = 9Z^2 e^2 q / 16 \text{ (with } q = 0.85 Z^{1/3} / a_H) = -E$$

Complete ionization (hyper-ionization, dissociation)  $|b_0| > a_H$

Critical field  $E > 7Z^4$ ; high field, for (X-ray) frequency  $\omega_0$

Partial ionization  $\delta Z$  ( $1 \ll \delta Z < Z$ ):  $E > 7(\delta Z)^4$

Ionization potential:  $\hbar\omega_0 = 30eV$  too large ( $6eV$ )

## Quasi-classical approximation

$$\dot{O} = \frac{i}{\hbar}[H, O], \quad \dot{O}_{mn} = \frac{i}{\hbar}(E_m - E_n)O_{mn} = i\omega_{mn}O_{mn}$$

Densely distributed states:  $\omega_{mn} \simeq -s(\partial E_m / \partial m)_m = -\omega_s$

Matrix elements  $O_{mn} = O_{m,m+s} \simeq O_s$  (temporal Fourier transform)

$$\dot{O} = -i\omega_s O + (\partial O^{cl} / \partial t)_{cl;h}; \quad O = O^{(1)} + iO^{(2)}, \quad \dot{O}^{(1)} = \omega_s O^{(2)} + (\partial O^{cl} / \partial t)_{cl;h}, \quad \dot{O}^{(2)} = -\omega_s O^{(1)}, \quad \ddot{O}^{(1)} + \omega_s^2 O^{(1)} = (\partial / \partial t)(\partial O^{cl} / \partial t)_{cl;h}$$

$$\ddot{O} + \omega_s^2 O = (\partial / \partial t)(\partial O^{cl} / \partial t)_{cl;h} = f$$

Classical harmonic oscillator; eigenfrequency  $\omega_s$  is the quantum transition frequency  $\omega_{mn}$  (wavepacket, lifetime  $\gamma^{-1}$ )

**QClass Ionization:** "Peripheral" electrons in heavy atoms are quasi-class

External field  $h = eEu \cos \omega t$ , force  $f = -eE \cos \omega t$

$$\ddot{u} + \omega_0^2 u + \gamma \dot{u} = -\frac{e}{m} E \cos \omega t$$

$\omega_0$  ( $= \omega_s$ ) - ionization frequency (exc states)

Amplitude  $|b_0| = eE/m\omega_0\gamma$ ;  $\gamma = 2r_0\omega_0^2/3c = 6 \times 10^{-24}\omega_0^2$

$\omega_0 = 10^{16}s^{-1}$  ( $\hbar\omega_0 = 6eV$ )  $\Rightarrow |b_0| = 10^{-7}E\text{ cm}$

Critical field  $E > 5 \times 10^{-2}\text{ statvolt/cm}$  ( $\simeq 1500V/m$ )

Higher fields: multiple-quanta transitions; Qclass: larger displacements, anharmonicities ( $\omega_0^2 = (1/m)(\partial^2 U / \partial u^2)_0$ )